

Section 1.

Dynamics (Newton's Laws of Motion)

Two approaches:

- 1) Given all the forces acting on a body, *predict* the subsequent (changes in) motion.
- 2) Given the (changes in) motion of a body, *infer* what forces act upon it.

Review of Newton's Laws:

First Law: A body at rest remains at rest,
a body in motion continues to move at constant velocity,
unless acted upon by an external force.

Note: This is only true in an inertial frame.

Example 1: You are on a train with a ball on the floor. The train accelerates. What does the ball do?

Example 2: You observe the world from a rotating carousel. All other objects are changing their velocities wrt you as you turn. What forces must be applied to them to achieve this?

Second Law: A force acting on a body causes an acceleration of the body,
in the direction of the force, proportional to the force,
and inversely proportional to the mass.

Note: This is only true in an inertial frame.

Express as $\vec{a} = \frac{\vec{F}}{m}$, or $\vec{F} = m\vec{a}$, or $\vec{F} = m \frac{d\vec{v}}{dt}$

Example 1: Force parallel to velocity. What does the body do?

Example 2: Force always at right angles to velocity. What does the body do?

We define the linear momentum $\vec{P} = m\vec{v}$, so that

$$\vec{F} = \frac{d\vec{P}}{dt}$$

We have two ways of *measuring* mass:

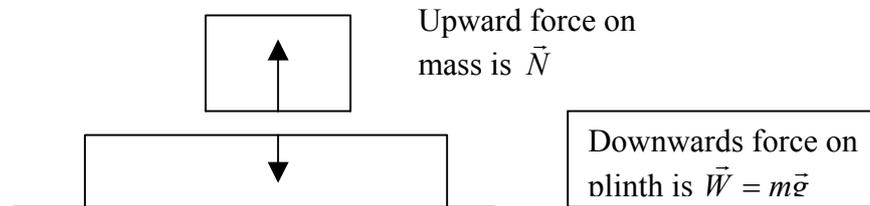
Inertial method – apply a force and measure acceleration $\Rightarrow m_{inertial}$

Gravitational method – weight the object (no motion) $\Rightarrow m_{grav}$

Later we will see that they are the same.

Third Law: To every action there is an equal and opposite reaction.

Example 1:



And $\vec{N} = -\vec{W}$

Example 2: A rocket engine generates the force \vec{F}_{thrust} and applies the force F_{gas} to the exhaust,

And $\vec{F}_{thrust} = -\vec{F}_{gas}$

Applications of Newton's Laws of Motion

Example 1: Inclined Ramp

An 8kg cart is pulled up a frictionless slope inclined at 20° . Determine the force if the cart is to move

- a) With uniform motion,
- b) With an acceleration of 0.2 m s^{-2} up the plane.

a) Resolve forces parallel to surface, then $F = 8g \sin 20^\circ = 26.8\text{N}$

b) Hence $F - 8g \sin 20^\circ = ma = 8 \times 0.2$, so $F = 28.4\text{N}$

Example 1: The Pulley

A weightless cord hangs over a frictionless pulley. A mass of 1kg hangs at one end of the cord and a mass of 2kg at the other. Calculate

- the acceleration of the masses,
- the tension in the cord,
- the reaction (the upwards force) exerted by the pulley.

Analysis: String must be at constant tension T throughout.

Let upwards acceleration be positive.

Let string accelerate at a on 1kg side

Then for the 1kg mass,

$$T - mg = ma, \text{ i.e. } T - g = a$$

and for the 2kg mass

$$T - mg = -ma \text{ i.e. } T - 2g = -2a$$

- Eliminating T , we obtain $a = g/3$
- From either equation, $T = 4/3 g$
- Then the reaction is $R = 2T = 8/3 g$

Now solve the same problem for masses m and m'

Equilibrium of a Solid Body

The static equilibrium of a solid body entails two distinct conditions:

- The net force tending to accelerate it is zero

$$\sum_i \vec{F}_i = 0 \quad \text{Condition of Translational Equilibrium}$$

Note: All the forces **do not** have to go through the same point. A ladder leaning against a wall has reaction and friction forces at each end, which do not go through the centre or any other single point.

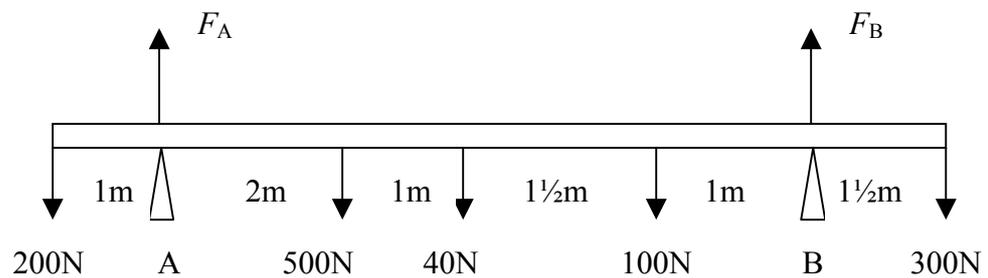
- The net torque tending to rotate it is zero

$$\sum_i \vec{T}_i = 0 \quad \text{Condition of Rotational Equilibrium}$$

Note: A body may be in translational equilibrium while out of rotational equilibrium. It may also be in rotational equilibrium while out of translational equilibrium.

Example: A Loaded Bar

A weightless bar rests on two supports. Several loads are hung from it



Calculate the forces F_A and F_B

Analysis:

- 1) Translational equilibrium, and note all forces are in y-direction. So forces simply sum to zero:

$$\sum_i F_i = 0 = 200 + 500 + 40 + 100 + 300 - F_A - F_B$$

$$F_A + F_B = 1140\text{N}$$

- 2) Rotational equilibrium, so total moment about any point is zero. Taking moments about A:

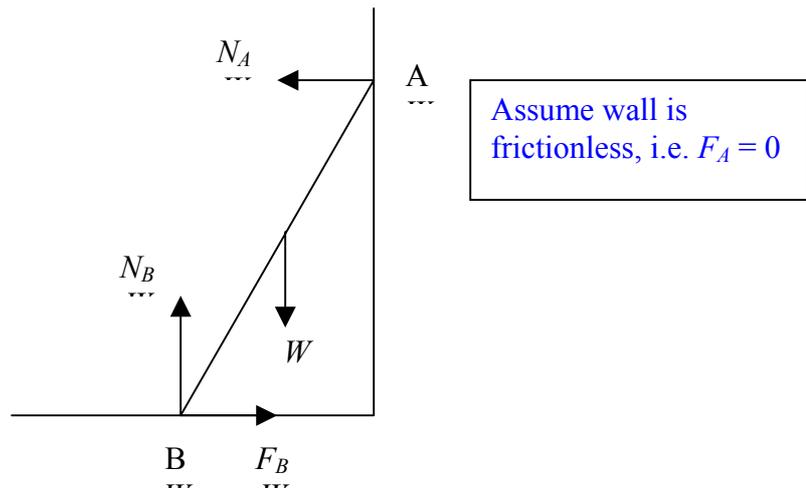
$$\sum_i T_i = \sum_i F_i x_i = 0$$

$$= 200 \times -1 - F_A \times 0 + 500 \times 2 + 40 \times 3 + 100 \times 4.5 - F_B \times 5.5 + 300 \times 7$$
$$= 3470 - 5.5F_B$$

$$\text{So } F_B = \frac{3470}{5.5} = 630.9\text{N}$$

$$\text{Therefore } F_A = 1140 - F_B = 509.1\text{N}$$

Example: Ladder against Wall



Let the weight of the ladder be 16 kg, so that $W = 160 \text{ N}$, and let it be at 60° ($\pi/3$ radians) to the ground.

Find the forces acting at each end of the ladder

Find the minimum coefficient of friction against the ground.

Analysis

Translation equilibrium: *Equate forces resolved in two axes*

$$F_B = N_A$$

$$W = N_B = \underline{\underline{160 \text{ N}}}$$

Rotational equilibrium: *Take moments about B*

$$W \times \frac{1}{2}L \cos 60^\circ = N_A \times L \sin 60^\circ$$

$$\Rightarrow N_A = 80 \cot 60^\circ = \underline{\underline{46.2 \text{ N}}}$$

Friction:

$$F_B = N_A \leq \mu_S N_B$$

$$\Rightarrow \mu_S \geq \underline{\underline{46.2 / 160 \approx 0.29}}$$

Frictional Forces

Friction is due to interactions between the atoms of an object and those of a surface that it touches. *Microscopic* roughness plays a role too. *Macroscopic* roughness is treated separately.

There are two kinds of frictional forces:

- *Static* Friction, when the surfaces are at rest
- *Kinetic* or Sliding Friction, when the surfaces are in relative motion.

In both cases, the frictional force *opposes* motion between the surfaces, and its magnitude is

$$F = \mu_S N$$

$$F = \mu_K N$$

$$\mu_S \geq \mu_K$$

where μ_S and μ_K are the *coefficients of static / kinetic friction* respectively.

Definitions:

- $F = \mu_S N$ is the *minimum* force required to set in motion the surfaces in contact and initially at rest, when the normal force (*contact force*) is N .
- $F = \mu_K N$ is the *minimum* force required to maintain the relative motion of the surfaces in contact, when the normal force (*contact force*) is N .

Example: Inclined Plane with Friction

Sliding Uphill: Resolving forces parallel to inclined plane,

and using $F = ma$,

$$F_{pull} - mg \sin \theta - \mu N = ma$$

a is uphill acceleration.

$$F_{pull} = mg \sin \theta + \mu mg \cos \theta + ma$$

Sliding Uphill: Resolving forces parallel to inclined plane,

and using $F = ma$,

$$F_{pull} + mg \sin \theta - \mu N = ma$$

a is downhill acceleration

$$F_{pull} = -mg \sin \theta + \mu mg \cos \theta +$$

$a = 0$ for
uniform motion.

Just balancing, $F_{pull} = 0 = a$, then $\mu = \tan \theta$

Acceleration with Varying Mass: The Rocket

A rocket at take-off has maximum mass (payload, structure and fuel) which decreases during flight as fuel is used. This is a characteristic example of a varying mass problem.

Let the rocket operate by ejecting exhaust at nozzle velocity v_e

at the rate $\frac{dm}{dt}$ (mass per unit time).

At time t let the rocket have velocity $v(t)$ relative to an inertial frame and mass $m(t)$. The exhaust mass dm that departs in the next time interval dt therefore has velocity $v'(t) = v(t) - v_e$ in the inertial frame.

Conserving momentum $p = mv$:

At time t , $p(t) = m(t)v(t)$

At time $t + dt$ $p(t + dt) = (m - dm)(v + dv) + v'dm$

which expands to $p(t + dt) = mv + mdv - vdm + v'dm - v_e dm$

The momentum is unchanged, so $mdv = v_e dm$

So acceleration is $a(t) = \frac{dv(t)}{dt} = \frac{v_e}{m(t)} \frac{dm(t)}{dt}$

Integrating this from initial conditions, at $t = 0$, $v = 0$, $m = m_0$,

$$v(t) = v_e \log_e \frac{m_0}{m(t)}$$

This is a very important equation for rocket designers. It says that the nozzle velocity must be made as high as possible, and that the fuel must be as high a proportion of the mass as possible. It implies that rockets should be multi-stage.

Work and Power

Work is done when a force acts along a displacement. The work is the energy required to achieve this.

Example: A body mass m falls through a height h . The work done by gravity is $Fh = mgh$.

Only the component of force parallel to the displacement is relevant. If the force is at the angle α to the displacement, the work is $Fs \cos \alpha$. In vector notation, this is

$$W = \vec{F} \cdot \vec{s}$$

If the path is curved, we may want the differential relationship,

$$dW = F \cos \alpha \, dr = \vec{F} \cdot d\vec{r}$$

Power is the rate of doing work.

$$P = \frac{dW}{dt} = \vec{F} \cdot \frac{d\vec{r}}{dt} = \vec{F} \cdot \vec{v}$$

Kinetic and Potential Energy

Consider $\vec{F} = m\vec{a}$ - a force acting on a particle to accelerate it. The work done is

$$\begin{aligned} W &= \int_A^B dW = \int_A^B \vec{F} \cdot d\vec{r} = m \int_A^B \vec{a} \cdot d\vec{r} = m \int_A^B \frac{d\vec{v}}{dt} \cdot d\vec{r} = m \int_A^B \frac{d\vec{r}}{dt} \cdot d\vec{v} \\ &= m \int_A^B \vec{v} \cdot d\vec{v} = \frac{1}{2}m(v_B^2 - v_A^2) \end{aligned}$$

We define *kinetic energy* accordingly as $\frac{1}{2}m\vec{v} \cdot \vec{v} = \frac{1}{2}mv^2$

Work done by a force accelerating a particle changes the kinetic energy of the particle.

Work may be done by or against a force due to a field. For example, a charged particle moving in an electric field under the Coulomb force, or a mass moving in a gravitational field.

In these cases we introduce the concept of *potential energy*. The work done on the particle by the force is

$$W_{ON} = \int_A^B \vec{F} \cdot d\vec{r} = \int_A^B dE = E_B - E_A$$

The work done by the particle is $W_{BY} = -W_{ON} = E_A - E_B$

In differential form,

$$dW = -\vec{F} \cdot d\vec{r} = -dE$$

Potential Energy Curves

We may write potential energy as a function of position, e.g. in one dimension, $E = E(x)$. Then

$$F(x) = -\frac{dE}{dx}$$

is the force acting on the particle. Consequently, *minima* and *maxima* of the potential energy curve are points where $F = 0$ and a stationary particle will remain at rest. The minima are stable (or metastable) and the maxima are unstable.
